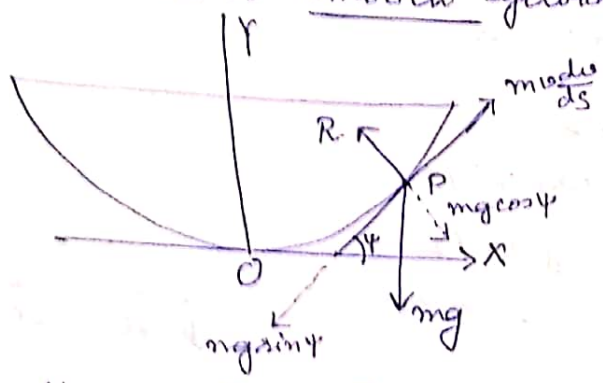


23/8/20

Motion in a Smooth cycloid under Gravity.



eqⁿ of cycloid,

$$s = 4a \sin \psi \quad \dots (1)$$

a = radius of generating circle.

$$\dot{s} = \frac{ds}{d\psi} = 4a \cos \psi$$

At any time P be the position of the particle.
eqⁿ of motion

$$m \dot{s} \frac{ds}{dt} = -mg \sin \psi \quad \dots (2)$$

$$m \frac{v^2}{\rho} = R - mg \cos \psi \quad \dots (3)$$

from (2) $\frac{d^2s}{dt^2} = -\frac{g}{4a} s \quad \dots (4)$

from (4) \rightarrow the motion of the cycloid is simple harmonic having a period $\frac{2\pi}{\sqrt{g/4a}} = 2\pi \sqrt{\frac{4a}{g}} = 4\pi \sqrt{\frac{a}{g}}$ which is independent of arc length.

from (1)

$$v \frac{dv}{ds} = -\frac{g}{4a} s$$

Integrating $\frac{v^2}{2} = -\frac{g}{4a} \frac{s^2}{2} + \frac{C}{2}$

$$v^2 = -\frac{g}{4a} s^2 + C$$

$$= -\frac{g}{4a} (4a \sin \psi)^2 + C$$

The particle starts from rest at $\psi = \psi_0$

$$0 = +\frac{g}{4a} C = 4ag \sin^2 \psi_0$$

$$v^2 = 4ag (\sin^2 \psi_0 - \sin^2 \psi)$$

$$R = mg \cos \psi + \frac{m}{4a \cos \psi} \cdot 4ag (\sin^2 \psi_0 - \sin^2 \psi)$$

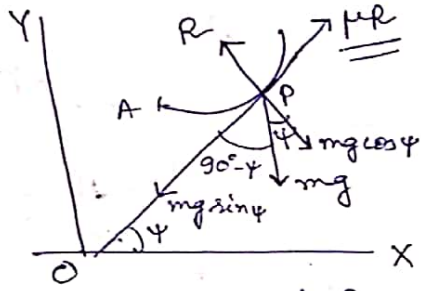
$$= mg \cos \psi + mg \sec \psi (\sin^2 \psi_0 - \sin^2 \psi)$$

$$= mg (\cos^2 \psi - \sin^2 \psi + \sin^2 \psi_0) \sec \psi$$

$$= mg (\cos 2\psi + \sin^2 \psi_0) \sec \psi$$

reaction at any pt. on the curve. Oscillation continues backwards & forwards. ~~passing through~~

Motion on a Rough Curve



Let a particle of mass m slide down a rough curve whose co-efficient of friction is μ . At any time t , P be the position of the particle. Distance of P from the fixed pt $A = s$.

R be the normal reaction at P

$$\text{Eqn of motion } m \frac{dv}{ds} = \mu R - mg \sin \psi \quad \dots (1)$$

$$\frac{mv^2}{\rho} = R - mg \cos \psi \quad \dots (2)$$

Eliminating R between (1) & (2), we have

$$v \frac{dv}{ds} - \frac{\mu v^2}{\rho} = g(\mu \cos \psi - \sin \psi)$$

$$\text{or, } \frac{1}{2} \frac{dv^2}{ds} - \frac{\mu v^2}{\rho} = g(\mu \cos \psi - \sin \psi)$$

$$\text{or, } \frac{1}{2} \cdot \frac{dv^2}{d\psi} \cdot \frac{d\psi}{ds} - \frac{\mu v^2}{\rho} = g(\mu \cos \psi - \sin \psi)$$

$$\text{or, } \frac{1}{2g} \left[\frac{dv^2}{d\psi} - 2\mu v^2 \right] = g(\mu \cos \psi - \sin \psi)$$

$$\text{or, } \frac{dv^2}{d\psi} - 2\mu v^2 = 2g(\mu \cos \psi - \sin \psi) \quad \dots (3)$$

$$\text{Int. factor } e^{-\int 2\mu d\psi} = e^{-2\mu\psi}$$

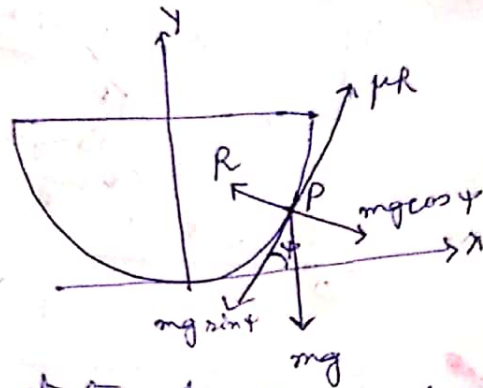
\therefore Multiplying both sides of (3) by $e^{-2\mu\psi}$ and integrating

$$v^2 e^{-2\mu\psi} = \int 2g(\mu \cos \psi - \sin \psi) e^{-2\mu\psi} d\psi + C \quad \dots (4)$$

For known curve ρ is known in terms of ψ .
 \therefore On integration v^2 , & from (2) R are obtained.

Motion in a Rough Cycloid under Gravity

Let a heavy particle of mass m slide inside a rough cycloid with μ its vertex at a lowest point and axis vertical. Let μ be the coefficient of friction on the curve.



Let the position of the particle at time t . Forces acting on the particle are shown in the fig.

Eqⁿ of motions

$$m \frac{dv}{ds} = \mu R - mg \sin \psi \quad \dots (1)$$

$$\frac{mv^2}{R} = R - mg \cos \psi \quad \dots (2)$$

Eliminating R between (1) & (2), we have

$$v \frac{dv}{ds} - \frac{v^2}{R} = g(\mu \cos \psi - \sin \psi)$$

$$\& \frac{1}{2} \frac{d(v^2)}{ds} - v^2 \frac{d\psi}{ds} = g(\mu \cos \psi - \sin \psi)$$

$$\& \frac{dv^2}{d\psi} - 2\mu v^2 = 2g(\mu \cos \psi - \sin \psi) \frac{ds}{d\psi} \quad \dots (3)$$

$$s = 4a \sin \psi \quad [\text{intrinsic eqⁿ of cycloid}] \quad \dots (4)$$

$$\frac{ds}{d\psi} = 4a \cos \psi$$

$$\frac{dv^2}{d\psi} - 2\mu v^2 = 8ag(\mu \cos \psi - \sin \psi) \cos \psi$$

$$\& \frac{d}{d\psi} (v^2 e^{-2\mu\psi}) = 4ag e^{-2\mu\psi} [\mu(1 + \cos 2\psi) - \sin 2\psi] \quad \dots (5)$$

$$\frac{d}{d\psi} (v^2 e^{-2\mu\psi}) = 4ag e^{-2\mu\psi} [\mu(1 + \cos 2\psi) - \sin 2\psi]$$

Integrating

$$v^2 e^{-2\mu\psi} = 4ag \int e^{-2\mu\psi} \{ \mu(1 + \cos 2\psi) - \sin 2\psi \} d\psi$$

If $v = 0$ at $\psi = \psi_0$ i.e., the particle starts from rest from a pt where the tangent makes an angle ψ_0 with the horizontal, then

$$v^2 e^{-2\mu\psi} = 4ag \int_{\psi_0}^{\psi} e^{-2\mu\psi} [\mu + \mu \cos 2\psi - \sin 2\psi] d\psi$$

Do yourself.

$$= 4ag \left[-\frac{e^{-2\mu\psi}}{2} + \mu e^{-2\mu\psi} \left\{ \frac{-2\mu \cos 2\psi + 2 \sin 2\psi}{4(1+\mu^2)} \right\} - e^{-2\mu\psi} \left\{ \frac{-2\mu \sin 2\psi - 2 \cos 2\psi}{4(1+\mu^2)} \right\} \right]_{\psi_0}^{\psi}$$

Simplifying

$$= 4ag \left[-\frac{e^{-2\mu\psi}}{2} + \frac{e^{-2\mu\psi}}{2(1+\mu^2)} \left\{ (1-\mu^2) \cos 2\psi + 2\mu \sin 2\psi \right\} \right]_{\psi_0}^{\psi}$$

$$= \frac{2ag}{1+\mu^2} \left[-e^{-2\mu\psi} \left\{ (1+\mu^2) + (1-\mu^2) \cos 2\psi - 2\mu \sin 2\psi \right\} \right]_{\psi_0}^{\psi}$$

$$= \frac{2ag}{1+\mu^2} \left[e^{-2\mu\psi} \left\{ \mu^2 (1 + \cos 2\psi) - (1 - \cos 2\psi) + 2\mu \sin 2\psi \right\} \right]_{\psi_0}^{\psi}$$

$$= \frac{2ag}{1+\mu^2} \left[e^{-2\mu\psi} \left\{ -2\mu^2 \cos^2 \psi - 2 \sin^2 \psi + 4\mu \sin \psi \cos \psi \right\} \right]_{\psi_0}^{\psi}$$

$$= \frac{2ag}{1+\mu^2} \cdot 2 \left[e^{-2\mu\psi} (\sin \psi - \mu \cos \psi)^2 \right]_{\psi_0}^{\psi}$$

$$= \frac{4ag}{1+\mu^2} \left[e^{-2\mu\psi_0} (\sin \psi_0 - \mu \cos \psi_0)^2 - e^{-2\mu\psi} (\sin \psi - \mu \cos \psi)^2 \right] \quad \text{--- (6)}$$

∴ using (6) R can be obtained from (2).

$$v = \frac{ds}{dt} = 4a \cos \psi \frac{d\psi}{dt}$$

$$\int_0^t dt = \int_{\psi_0}^{\psi} \frac{4a \cos \psi}{v} d\psi$$

v is a fⁿ of ψ.

we can integrate for the given problem.

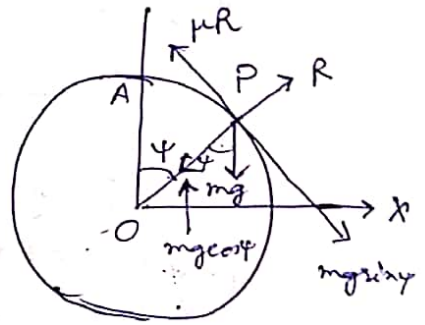
Motion on a Rough Circle under Gravity.

P be the position of the particle at any time t .

$$\overline{AP} = s, \quad s = a\psi$$

R is the normal reaction

μR = tangential frictional force acting opposite to the motion.



Eqs of motion: -

$$m v \frac{dv}{ds} = mg \sin \psi - \mu R \quad \text{--- (1)}$$

$$s = a\psi \\ \frac{ds}{d\psi} = a$$

$$\frac{m v^2}{s} = R - mg \cos \psi \quad \text{--- (2)}$$

Eliminating R from the above equations

$$m v \frac{dv}{d\psi} \cdot \frac{d\psi}{ds} + \frac{m v^2}{s} \mu = mg \sin \psi - mg \cos \psi$$

$$\Rightarrow 2 v \frac{dv}{d\psi} + 2 \mu v^2 = 2ga (\sin \psi - \mu \cos \psi)$$

$$\Rightarrow \frac{dv^2}{d\psi} + 2 \mu v^2 = 2ga (\sin \psi - \mu \cos \psi)$$

Int. fact. $\frac{d}{d\psi} (v^2 e^{2\mu\psi}) = 2ga e^{2\mu\psi} (\sin \psi - \mu \cos \psi)$

$$\text{Int } v^2 e^{2\mu\psi} = 2ga e^{2\mu\psi} \frac{2\mu \sin \psi - \cos \psi - 2\mu^2 \cos \psi + \mu \sin \psi}{1+4\mu^2} + C$$

$$v^2 = \frac{2ga}{1+4\mu^2} \cdot (\mu \sin \psi - (1+2\mu^2) \cos \psi) + C e^{-2\mu\psi} \quad \text{--- (3)}$$

$$\int e^{ax} \sin bx dx = e^{ax} \frac{a \sin bx - b \cos bx}{a^2 + b^2}$$

$$\int e^{ax} \cos bx dx = e^{ax} \frac{a \cos bx + b \sin bx}{a^2 + b^2}$$

at $\psi = 0, \quad v = v_0$

$$v_0^2 = C - \frac{2ag}{1+4\mu^2} (1+2\mu^2)$$

$$\therefore C = v_0^2 + \frac{2ag}{1+4\mu^2} (1+2\mu^2) \quad \text{--- (4)}$$

Substituting in (4) in (3), v^2 is obtained,
 $\&$ from (2) using v^2 - - R can be obtained.

Ex A particle slides down the smooth curve $y = a \sinh \frac{x}{a}$, the axis of x being horizontal and the axis of y downwards, starting from rest at the point where the tangent is inclined at α to the horizon; show that it will leave the curve when it has fallen through a vertical distance $a \sec \alpha$.

$$y = a \sinh \frac{x}{a}$$

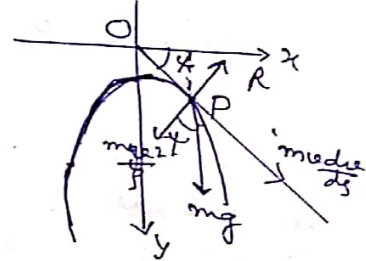
P be the position of the particle at time t .

$R \rightarrow$ normal reaction -

Eqⁿ of motion -

$$m \frac{dv}{ds} = mg \sin \psi \quad \dots (1)$$

$$\frac{mv^2}{\rho} = mg \cos \psi - R \quad \dots (2)$$



$$y = a \sinh \frac{x}{a}$$

$$\frac{dy}{dx} = a \cdot \frac{1}{a} \cosh \frac{x}{a}$$

$$y_1 = \tan \psi = \cosh \frac{x}{a}$$

$$= \sqrt{1 + \sinh^2 \frac{x}{a}}$$

$$= \sqrt{1 + \frac{y^2}{a^2}}$$

$$\frac{d^2y}{dx^2} = \frac{d}{dx} \cosh \frac{x}{a} = \frac{1}{a} \sinh \frac{x}{a} = \frac{y}{a^2}$$

$$\rho = \frac{(1 + y_1^2)^{3/2}}{y_2} = \frac{(1 + \cosh^2 \frac{x}{a})^{3/2}}{1/a}$$

$$= \frac{(1 + 1 + \frac{y^2}{a^2})^{3/2}}{1/a}$$

$$= \frac{(\frac{y}{a^2})^{3/2} (2a^2 + y^2)^{3/2}}{\frac{y}{a^2} a^3} = \frac{(2a^2 + y^2)^{3/2}}{ay}$$

When $\psi = \alpha$

$$y = b$$

$$\tan \alpha = \sqrt{1 + \frac{b^2}{a^2}}$$

from (1)

$$v \frac{dv}{ds} = g \sin \psi$$

$$= g \cdot \frac{dy}{ds}$$

$$\int 2v dv = 2g \int dy$$

$$\text{Int } v^2 = 2gy + C$$

$$v = 0, \text{ at } y = b, \quad C = -2gb$$

$$v^2 = 2g(y - b) \quad \dots (3)$$

The particle will leave the plane when $R = 0$.

$$2g(y - b) = g \rho \cos \psi$$

$$\int 2g(y - b) = g \rho \cos \psi = \frac{g \rho}{\sqrt{1 + \tan^2 \psi}}$$

$$\int 2(y - b) = \frac{g(2a^2 + y^2)^{3/2}}{ay \sqrt{1 + \frac{y^2}{a^2}}} = \frac{g(2a^2 + y^2)^{3/2}}{ay \frac{\sqrt{2a^2 + y^2}}{a}} = \frac{g(2a^2 + y^2)}{y}$$

$$\int 2(y - b) = g \cdot \frac{2a^2}{y} + y$$

$$\int y - 2b = \frac{2a^2}{y}, \quad \text{or}$$

3) $y^2 - 2by = 2a^2$

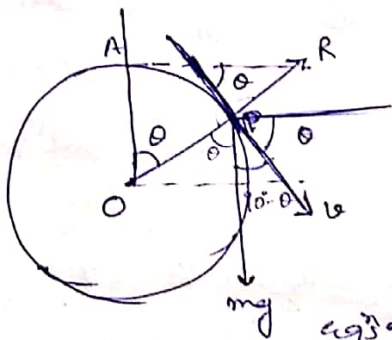
3) $y^2 - 2by + b^2 = 2a^2 + b^2$

3) $(y-b)^2 = a^2 + a^2 \tan^2 \alpha$
 $= a^2 \sec^2 \alpha$

$\therefore y - b = a \sec \alpha$, proved.

$\sec \alpha = 1 + \frac{b^2}{a^2}$
 $a^2 \tan^2 \alpha = a^2 + b^2$

Ex A particle starts from rest at the highest point on the outside of a smooth circle. Show that the path described after leaving the circle is the parabola, whose latus rectum of the parabola, which the particle describes after leaving the circle is $\frac{16}{27}$ times the radius of the circle.



O - centre of the circle
 P - position of the particle at any time t.
 $\angle AOP = \theta$
 $AP = s$
 v - vel of the particle,
 R - normal reaction

eqn of motion
 $m \frac{dv}{dt} = mg \sin \theta$ (1)

$s = at$
 $\frac{ds}{dt} = a$

$\frac{mv^2}{s} = mg \cos \theta - R$ (2)

from (1) $v \frac{dv}{d\theta} = g \sin \theta$

or $v dv = g \sin \theta d\theta$

Int $v^2 = -2g \cos \theta + C$

Initially $t = 0, \theta = 0, v = 0 \therefore C = 2ag$

$v^2 = 2ag(1 - \cos \theta)$ (3)

from eqn (2) $R = mg \cos \theta - \frac{mv^2}{s}$
 $= mg \cos \theta - \frac{2mg(1 - \cos \theta)}{2a \cos \theta} (1 - \cos \theta)$
 $= mg(3 \cos \theta - 2)$ (4)

$R = 0 \rightarrow$ the particle leaves the circle, when $\theta = \alpha$
 $3 \cos \alpha - 2 = 0, \cos \alpha = \frac{2}{3}$

$v^2 = 2ga(1 - \frac{2}{3}) = \frac{2}{3} ag$

$v = \sqrt{\frac{2ag}{3}}$

horizontal component is $v \cos \alpha$
 $= \sqrt{\frac{2}{3}ag} \cos \alpha$. $\cos \alpha = \frac{2}{3}$
 $= \frac{2}{3} \sqrt{\frac{2}{3}ag}$

After leaving the circle, the particle describes parabolic path with vel. $\frac{2}{3} \sqrt{\frac{2}{3}ag}$.

Latus rectum of the parabola = $\frac{2}{g} \times (\text{hor. comp. of vel.})^2$
 $= \frac{2}{g} \cdot \frac{4}{9} \cdot \frac{2}{3} ag = \frac{16a}{27}$ proved.

Chosh
 & Chakraborty

(18)

A particle is projected along the inner surface of a rough sphere and is acted upon by no forces. Show that it will return to the point of projection after time $\frac{a}{\mu v} (e^{2\mu} - 1)$, where a is the radius of the sphere, μ is the coefficient of friction and v is the velocity of projection.



The particle is projected from the lowest Pt A in the sphere. P be the position at time t.

$\widehat{AP} = s$
 eqⁿ of motion $m \frac{dv}{ds} = -\mu R$ --- (1)

$\frac{mv^2}{R} = \mu R$ --- (2)

$\therefore -\frac{v dv}{ds} = -\frac{\mu v^2}{R}$ $s = a\theta$
 $\frac{ds}{d\theta} = a = R$

$\frac{dv^2}{ds} + \frac{2\mu v^2}{R} = 0$

$\frac{dv^2}{d\theta} \cdot \frac{d\theta}{ds} + \frac{2\mu}{a} v^2 = 0$

$\frac{dv^2}{d\theta} + 2\mu v^2 = 0$

$\frac{d}{d\theta} (v^2 e^{2\mu\theta}) = 0$

$\int v^2 e^{2\mu\theta} = C$

Initially $v = V, \theta = 0, C = V^2$

$\therefore v^2 e^{2\mu\theta} = V^2$

$\therefore v e^{\mu\theta} = V$

Q

$$07. \quad \frac{ds}{dt} e^{\mu\theta} = v$$

$$\& \quad \frac{ds}{d\theta} \cdot \frac{d\theta}{dt} \cdot e^{\mu\theta} = v$$

$$\theta = 0 \text{ when } t = 0$$

$$\theta = 2\pi \text{ when } t = t_1$$

$$\& \quad a e^{\mu\theta} d\theta = v dt$$

$$\& \quad \int_0^{2\pi} a e^{\mu\theta} d\theta = \int_0^{t_1} v dt$$

$$t_1 = \frac{a}{\mu v} [e^{\mu\theta}]_0^{2\pi}$$

$$= \frac{a}{\mu v} [e^{2\mu\pi} - 1]. \text{ proved.}$$

Ghosh Chakr. (22)
Saha Ganguly (33)

A particle slides from a cusp down the arc of a rough cycloid, the axis of which is vertical. Prove that its velocity at the point, when the cycloid is smooth, the ratio of $(e^{-\mu\pi} - \mu^2)^{\frac{1}{2}} : (1 + \mu^2)^{\frac{1}{2}}$, where μ is the coefficient of friction.

Let the position of the particle

eqn of motion

$$\frac{m v^2}{ds} = \mu R - mg \sin \psi$$

$$\frac{m v^2}{s} = R - mg \cos \psi$$

$$\frac{1}{2} \frac{dv^2}{ds} - \frac{\mu v^2}{s} = g (\mu \cos \psi - \sin \psi)$$

$$\& \quad \frac{dv^2}{d\psi} \cdot \frac{d\psi}{ds} - \frac{2\mu v^2}{s} = 2g (\mu \cos \psi - \sin \psi)$$

$$\& \quad \frac{dv^2}{d\psi} - 2\mu v^2 = 2sg (\mu \cos \psi - \sin \psi)$$

$$\& \quad \frac{d}{d\psi} [v^2 e^{-2\mu\psi}] = 2g \int s e^{-2\mu\psi} (\mu \cos \psi - \sin \psi) d\psi$$

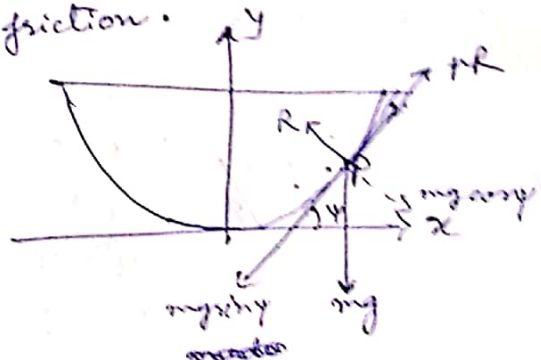
I.F. $e^{-2\mu\psi}$, $s = 4a \sin \psi$
 $\frac{ds}{d\psi} = 4a \cos \psi$

$$\int v^2 e^{-2\mu\psi} = 8ag \int e^{-2\mu\psi} (\mu \cos \psi - \sin \psi) \cos \psi d\psi + A$$

Let $z = e^{-\mu\psi} (\mu \cos \psi - \sin \psi)$

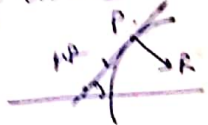
$$dz = [e^{-\mu\psi} (\mu \sin \psi - \cos \psi) - \mu e^{-\mu\psi} (\mu \cos \psi - \sin \psi)] d\psi$$

$$= -e^{-\mu\psi} \cos \psi (1 + \mu^2) d\psi$$



11.10
 1. A heavy bead slides on a smooth fixed circular wire of radius a . It is projected from the lowest point with a velocity just sufficient to carry it to the highest pt. Prove that in time t the radius through the bead will turn through an angle $2 \tan^{-1}(\sinh \sqrt{\frac{g}{a}} t)$, and that the bead will take an infinite time to reach the highest point.
 [hint: same as ~~part~~ ~~on~~ ~~the~~ ~~(~~above~~)~~ ~~].~~ ~~2020~~].

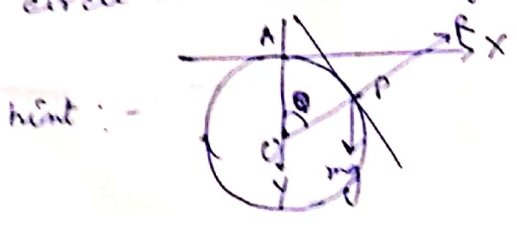
2. A bead moves along a rough curved wire which is such that it changes its direction of motion with constant angular velocity. Such that a possible form of the wire is an equiangular spiral.

[hint:  $m \frac{dv}{ds} = -R$ $\frac{d\psi}{dt} = \omega \text{ const}$]

$\frac{mv^2}{S} = R$

eqⁿ of equiangular spiral $S = ae^{k\psi}$

3. A particle starts from rest and slides from the highest point of a smooth circle. Prove that it will leave the curve when it has fallen a vertical height of one-third the radius of the circle and its velocity in then is $\sqrt{\frac{2ag}{3}}$.



Pr Deter - submit. Submit on 27.8.2020 8-30 a.m.

If v_1 be the velocity at the lowest point ($y = 0$), then

$$\int_0^{v_1} v \, dv = -g \int_{20}^0 dy, \quad \text{or,} \quad \frac{1}{2}v_1^2 = -g(0 - 20) = 20g$$

$$\text{or,} \quad v_1^2 = 40g; \quad \therefore \quad v_1 = \sqrt{40g} = 2\sqrt{10g}.$$

Normal equation of motion is $m\frac{v^2}{\rho} = R - mg \cos \psi$, where R is the normal pressure.

$$\text{or,} \quad R = m \left(\frac{v^2}{\rho} + g \cos \psi \right)$$

$$\text{and} \quad \rho = \frac{(1 + y_1^2)^{3/2}}{y_2} = \frac{\left(1 + \frac{4x^2}{25}\right)^{3/2}}{2/5} = \frac{5}{2} \cdot \frac{(25 + 4x^2)^{3/2}}{125}.$$

From (2),

$$v^2 = 2g(20 - y) = 2g \left(20 - \frac{1}{5}x^2\right) = \frac{2g}{5}(100 - x^2).$$

$$\therefore R = m \cdot \left\{ \frac{2g}{5}(100 - x^2) \times \frac{50}{(25 + 4x^2)^{3/2}} + g \cdot \frac{5}{\sqrt{4x^2 + 25}} \right\}$$

$$= 5mg \left\{ \frac{400 - 4x^2 + 25 + 4x^2}{(25 + 4x^2)^{3/2}} \right\} = 5mg \cdot \frac{425}{(25 + 4x^2)^{3/2}}.$$

At $x = 5$,

$$R = 5mg \cdot \frac{425}{(25 + 100)^{3/2}} = 5mg \cdot \frac{425}{125(5)^{3/2}} = \frac{17}{5\sqrt{5}}mg.$$

□ EXAMPLE 11. A particle is projected with velocity V from the cusp of a smooth inverted cycloid down the arc; show that the time of reaching the vertex is

$$2\sqrt{\frac{a}{g}} \tan^{-1} \left[\sqrt{\frac{4ag}{V}} \right].$$

[C.U. B.A./B.Sc.(H)'79]

SOLUTION. Let P be the position of the particle at time t on the smooth cycloid AOB , where s is the arcual distance of P from the vertex O .

Let R be the normal reaction at P .

The equations of motion are

$$mv \frac{dv}{ds} = -mg \sin \psi \quad (1)$$

$$\text{and } m \frac{v^2}{\rho} = R - mg \cos \psi, \quad (2)$$

where mg is the weight of the particle acting vertically downwards.

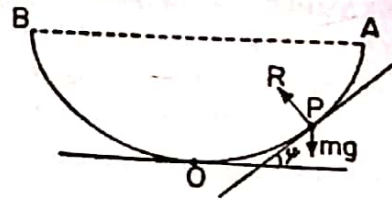


Fig. 7.20

Now, the equation of the curve is

$$s = 4a \sin \psi, \quad \text{or, } \sin \psi = s/4a;$$

\therefore From (1), $v \frac{dv}{ds} = -\frac{gs}{4a}$ and integrating, $v^2 = -\frac{g}{4a} \cdot s^2 + C$, where C is the constant of integration.

Initially, when $s = 4a, v = -V$; $\therefore C = V^2 + 4ag$.

$$\text{Hence, } v^2 = \frac{g}{4a}(16a^2 - s^2) + V^2;$$

$$\therefore \frac{ds}{dt} = v = -\sqrt{\frac{g}{4a}(16a^2 - s^2) + V^2},$$

[the negative sign is taken, since the particle is moving down the curve, i.e., in the direction of s decreasing.]

$$\text{or, } \sqrt{\frac{g}{4a}} dt = -\frac{ds}{\sqrt{16a^2 + \frac{4aV^2}{g} - s^2}}$$

Integrating,

$$\frac{1}{2} \sqrt{\frac{g}{a}} \int_0^T dt = - \int_{4a}^0 \frac{ds}{\sqrt{16a^2 + \frac{4aV^2}{g} - s^2}}$$

[since, when $t = 0, s = 4a$ and when $t = T$ (say), $s = 0$, i.e., the particle reaches the vertex O at time T .]

$$\begin{aligned} \therefore \frac{1}{2} \sqrt{\frac{g}{a}} T &= \left[\cos^{-1} \frac{s}{\sqrt{16a^2 + \frac{4aV^2}{g}}} \right]_{4a}^0 = \frac{\pi}{2} - \cos^{-1} \frac{4a}{\sqrt{16a^2 + \frac{4aV^2}{g}}} \\ &= \sin^{-1} \frac{4a}{\sqrt{16a^2 + \frac{4aV^2}{g}}} = \sin^{-1} \frac{1}{\sqrt{1 + \frac{V^2}{4ag}}} = \tan^{-1} \frac{\sqrt{4ag}}{V}, \end{aligned}$$

$$\text{or, } T = 2 \sqrt{\frac{a}{g}} \tan^{-1} \frac{\sqrt{4ag}}{V}.$$

EXAMPLE 12. A cycloid is placed with its axis vertical and vertex upwards and a heavy particle is projected from the cusp up the concave side of the curve with velocity $\sqrt{2gh}$; prove that the latus rectum of a parabola described after leaving the arc is $\frac{h}{2a}$ where a is the radius of the generating circle.

SOLUTION. Let the particle be projected from the cusp A where $s = 4a$ and let P be the position of the particle at any time t , where arc $OP = s$.

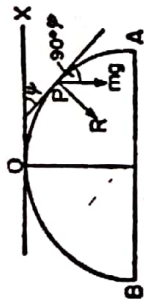


Fig. 7.21

From (1),

$$\frac{dv}{ds} = \frac{gs}{4a}, \text{ or, } v dv = \frac{g}{4a} s ds.$$

Integrating, $v^2 = \frac{gs^2}{4a} + A_1$, where A_1 is the constant of integration. Initially at A , $v = \sqrt{2gh}$, $s = 4a$; $\therefore 2gh = 4ag + A_1$.

$$\therefore v^2 = \frac{gs^2}{4a} + 2gh - 4ag.$$

Putting $s = 4a \sin \psi$,

$$v^2 = 2gh + 4ag \sin^2 \psi - 4ag$$

$$\text{or, } v^2 = 2gh - 4ag \cos^2 \psi.$$

From (2), using (3), we get

$$R = \frac{m(2gh - 4ag \cos^2 \psi)}{4a \cos \psi} - mg \cos \psi \left[\therefore \rho = \frac{ds}{d\psi} = 4a \cos \psi \right]$$

$$= \frac{mg}{4a \cos \psi} [2h - 8a \cos^2 \psi].$$

The particle will leave the cycloid, when $R = 0$, i.e., when $2h - 8a \cos^2 \psi = 0$, or, $\cos^2 \psi = \frac{h}{4a}$, which gives the value of ψ when it leaves the arc.

If u be the velocity with which the particle leaves the arc, then from (3), using $\cos^2 \psi = \frac{h}{4a}$, we get

$$u^2 = 2gh - 4ag \cdot \frac{h}{4a} = gh.$$

Now, we know that the latus rectum of the parabolic path is $\frac{2u^2 \cos^2 \alpha}{g}$, when a particle is projected with velocity u and at an angle α .

Hence the latus rectum of the parabola described after leaving the arc

$$= \frac{2u^2 \cos^2 \alpha}{g} = \frac{2u^2 \cos^2 \psi}{g} = \frac{2gh \cdot \frac{h}{4a}}{g} = \frac{h^2}{2a}.$$

EXAMPLE 13. A bead slides on a smooth parabolic wire held in a vertical plane with axis vertical and vertex downwards starting from rest at a height h above the vertex. Show that the square of the normal reaction at any point is inversely proportional to the cube of the height of the point above the directrix.

[C.U. B.A./B.Sc.(H) '52, '64]

SOLUTION. Let $P(x, y)$ be the position of the bead of mass m at any time t with reference to the vertex O as origin, the tangent Ox at O as the x -axis and the upward vertical through O as the y -axis.

Then the equation of the parabola is

$$x^2 = 4ay. \quad (1)$$

The equations of motion are

$$m \frac{dv}{ds} = -mg \cos(90^\circ - \psi) \quad (2)$$

$$\text{and } m \cdot \frac{v^2}{\rho} = R - mg \sin(90^\circ - \psi). \quad (3)$$

From (2),

$$\frac{dv}{ds} = -g \sin \psi = -g \frac{dy}{ds}, \text{ or, } 2v dv = -2g dy;$$

$$\therefore v^2 = -2gy + A,$$

where A is the constant of integration.

Initially, at $y = h$, $v = 0$; $\therefore 0 = -2gh + A$, or, $A = 2gh$.

$$\therefore v^2 = 2g(h - y).$$

From (3),

$$R = \frac{mv^2}{\rho} + mg \cos \psi = \frac{2mg(h - y)}{\rho} + mg \cos \psi.$$

$$= \frac{mg}{\rho} [2(h - y) + 2\rho \cos \psi]. \quad (4)$$

Now, from (1),

$$\frac{dy}{dx} = \frac{x}{2a} = \tan \psi$$

$$\therefore x = 2a \tan \psi \quad (5)$$

$$\text{or, } \frac{dx}{d\psi} = 2a \sec^2 \psi = 2a \left(1 + \frac{x^2}{4a^2} \right) = 2a \left(1 + \frac{4ay}{4a^2} \right)$$

[from (5) and (1)]

$$\text{or, } \frac{dx}{ds} \cdot \frac{ds}{d\psi} = 2a \left(1 + \frac{y}{a} \right) = 2(a + y)$$

$$\text{or, } \rho \cos \psi = 2(a + y). \quad (6)$$

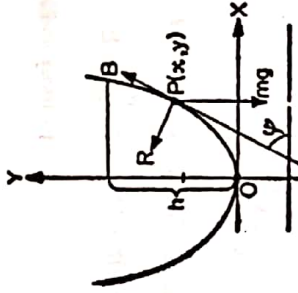


Fig. 7.22

$$\text{or, } 2v \cdot \frac{dv}{d\psi} - 2\mu v^2 = 2g(\mu \cos \psi - \sin \psi) 4a \cos \psi$$

[Since for a cycloidal curve, $s = 4a \sin \psi$

$$\text{or, } \frac{d(v^2)}{d\psi} - 2\mu \cdot v^2 = 4ag(\mu \cdot 2 \cos^2 \psi - 2 \sin \psi \cos \psi)$$

$$\text{or, } \frac{d(v^2)}{d\psi} - 2\mu \cdot v^2 = 4ag\{\mu(1 + \cos 2\psi) - \sin 2\psi\}. \quad (3)$$

This is a linear equation in v^2 and ψ , whose integrating factor is

$$e^{\int -2\mu d\psi}, \text{ i.e., } e^{-2\mu\psi}$$

Hence, the solution of (3) is

$$v^2 \cdot e^{-2\mu\psi} = 4ag \int e^{-2\mu\psi} \{\mu(1 + \cos 2\psi) - \sin 2\psi\} d\psi + C,$$

where C is the constant of integration,

$$\begin{aligned} \text{or, } v^2 e^{-2\mu\psi} &= 4ag \left[-\frac{e^{-2\mu\psi}}{2} + \frac{\mu e^{-2\mu\psi}}{4(1 + \mu^2)} (-2\mu \cos 2\psi + 2 \sin 2\psi) \right. \\ &\quad \left. - \frac{e^{-2\mu\psi}}{4(1 + \mu^2)} (-2\mu \sin 2\psi - 2 \cos 2\psi) \right] + C \\ &= \frac{ag e^{-2\mu\psi}}{1 + \mu^2} \left[-2(1 + \mu^2) + \mu(-2\mu \cos 2\psi + 2 \sin 2\psi) \right. \\ &\quad \left. + (2\mu \sin 2\psi + 2 \cos 2\psi) \right] + C \\ &= \frac{2ag \cdot e^{-2\mu\psi}}{1 + \mu^2} \left[-(1 + \mu^2) + (1 - \mu^2) \cos 2\psi + 2\mu \sin 2\psi \right] + C \\ &= -\frac{2ag \cdot e^{-2\mu\psi}}{1 + \mu^2} [2 \sin^2 \psi + 2\mu^2 \cos^2 \psi - 4\mu \sin \psi \cos \psi] + C \\ &= -\frac{4ag \cdot e^{-2\mu\psi}}{1 + \mu^2} (\sin \psi - \mu \cos \psi)^2 + C. \end{aligned} \quad (4)$$

Now, since $v = 0$ when $\psi = 0$ and $\frac{\pi}{2}$, from (4) we have

$$0 = -\frac{4ag}{1 + \mu^2} \cdot \mu^2 + C \quad (5)$$

$$\text{and } 0 = -\frac{4ag \cdot e^{-\mu\pi}}{1 + \mu^2} + C. \quad (6)$$

From (5) and (6), we get $\mu^2 = e^{-\mu\pi}$, or, $\mu^2 e^{\mu\pi} = 1$.

EXAMPLE 16. A particle is projected along the inner surface of a rough sphere and is acted on by no forces. Show that it will return to the point of projection at the end of time $\frac{2a}{g} (e^{2\pi\mu} - 1)$, where a is the radius of the sphere, V is the velocity of projection and μ is the coefficient of friction.

[B.U.(H)92; C.U.(H)78; 91; V.U.(H)91; SOLUTION. Since the particle is acted on by no forces, it will continue to move in a circle (circular section of the sphere).

Let the particle be projected from the point A with velocity V and P be its position at any time t such that

$$\angle AOP = \theta \text{ and arc } AP = s.$$

Then $s = a\theta$, $\dot{s} = a\dot{\theta}$ and $\ddot{s} = a\ddot{\theta}$.

If R be the normal reaction at P , then μR is the force of friction acting opposite to the direction of motion. The equations of motion are

$$\begin{aligned} \ddot{s} &= -\mu R \text{ and } \frac{v^2}{\rho} = R, \\ \text{or, } a\ddot{\theta} &= -\mu R \quad (1) \\ \text{and } a\dot{\theta}^2 &= R \quad [\because \rho = a] \quad (2) \end{aligned}$$

From (1) and (2), eliminating R , we get,

$$a\ddot{\theta} = -\mu a\dot{\theta}^2, \text{ or, } \ddot{\theta} + \mu\dot{\theta}^2 = 0, \text{ or, } \frac{\ddot{\theta}}{\dot{\theta}} + \mu\dot{\theta} = 0.$$

Integrating, $\log \dot{\theta} + \mu\theta = C$, where C is a constant.

Initially, at A , $\theta = 0, \dot{\theta} = \frac{V}{a} = \frac{V}{a}$; $\therefore \log \frac{V}{a} = C$.

$$\therefore \log \dot{\theta} + \mu\theta = \log \frac{V}{a}, \text{ or, } \log \left(\frac{\dot{\theta}}{V/a} \right) = -\mu\theta;$$

$$\therefore \dot{\theta} = \frac{V}{a} \cdot e^{-\mu\theta}, \text{ or, } dt = \frac{a}{V} e^{\mu\theta} d\theta.$$

The particle will return to the point of projection when $\theta = 2\pi$. If T be the required time, then

$$\int_0^T dt = \int_0^{2\pi} \frac{a}{V} e^{\mu\theta} \cdot d\theta, \text{ or, } T = \frac{a}{V} \left[\frac{e^{\mu\theta}}{\mu} \right]_0^{2\pi}$$

$$\text{Hence } T = \frac{a}{\mu V} (e^{2\pi\mu} - 1).$$

EXAMPLE 17. A particle is projected horizontally from the lowest point of a rough sphere of radius a , and comes to rest at the lowest point after describing an arc less than a quadrant, show that the velocity of projection must be $\sin \alpha \sqrt{2ga(1 + \mu^2)} / (1 - 2\mu^2)$, where α is the angle subtended by the arc at the centre and μ is the coefficient of friction.

[C.U. B.A./B.Sc.(H)70; 95; (P)85]

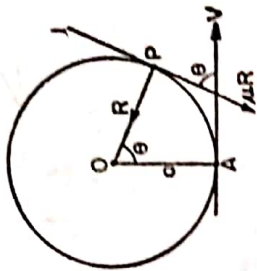


Fig. 7.25

Ex. 2. A smooth parabolic tube is placed, vertex downwards, in a vertical plane. A particle slides down the tube from rest under the influence of gravity. Prove that, in any position, the reaction of the tube is $\frac{2w(h+a)}{\rho}$, where w is the weight of the particle, ρ is the radius of curvature, $4a$ is the latus rectum and h is the original height of the particle above the vertex. [C. H. 1996]

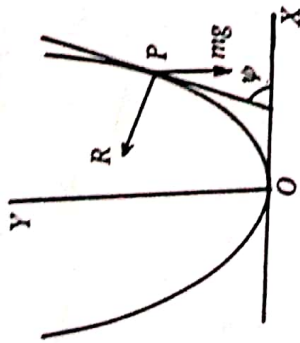


Fig. 13.8 (a)

Let the equation of the parabola with axis as the y-axis be $x^2 = 4ay$.

$$\text{Therefore } \rho = \frac{\left\{1 + \left(\frac{dy}{dx}\right)^2\right\}^{\frac{3}{2}}}{d^2y/dx^2}$$

$$\text{or, } \rho = 2a(1 + \tan^2 \psi)^{\frac{3}{2}}, \text{ since } \frac{dy}{dx} = \tan \psi = \frac{x}{2a}$$

$$= 2a \sec^3 \psi$$

$$= \frac{2a}{\cos \psi} \left(1 + \frac{x^2}{4a^2}\right) = \frac{2a}{\cos \psi} \left(1 + \frac{y}{a}\right) = \frac{2(a+y)}{\cos \psi} \quad \dots (1)$$

Equation of motion along the tangent is

$$mv \frac{dv}{ds} = -mg \sin \psi$$

$$\text{or, } v \frac{dv}{ds} = -g \frac{dy}{ds}$$

Integrating, we get $v^2 = C - 2gy$.

When $y = h$, we have $v = 0$.

$$\text{Therefore } C = 2gh.$$

... (2)

Therefore $v^2 = 2g(h-y)$.
Equation of motion along the normal is

$$m \frac{v^2}{\rho} = R - mg \cos \psi.$$

$$\begin{aligned} \text{Therefore } R &= m \left(g \cos \psi + \frac{v^2}{\rho} \right) \\ &= m \left\{ g \cdot \frac{2(a+y)}{\rho} + \frac{1}{\rho} 2g(h-y) \right\}, \text{ by (1) and (2)} \\ &= 2mg \cdot \frac{h+a}{\rho} = \frac{2w(h+a)}{\rho}. \end{aligned}$$

Ex. 3. A particle falls down a cycloid under its own weight starting from the cusp. Show that when it arrives at the vertex the pressure on the curve is twice the weight of the particle.

We have, as usual, the equations of motion

$$v \frac{dv}{ds} = -g \sin \psi = -\frac{g}{4a} s, \quad \dots (1)$$

$$m \frac{v^2}{\rho} = R - mg \cos \psi. \quad \dots (2)$$

Integrating equation (1), we get

$$v^2 = -\frac{g}{4a} s^2 + C. \quad \dots (3)$$

Initially, $v = 0$, when $s = 4a$.

$$\text{Therefore } C = \frac{g}{4a} \cdot 16a^2.$$

$$\text{Hence (3) reduces to } v^2 = \frac{g}{4a} (16a^2 - s^2).$$

At the vertex, $s = 0$; hence $v^2 = 4ag$ at the vertex.

$$\text{Hence, from (2), } R = \frac{mv^2}{\rho} + mg \cos \psi = m \cdot \frac{4ag}{4a} + mg = 2mg,$$

since $s = 4a \sin \psi$ and $\rho = 4a \cos \psi$, so that, at $\psi = 0$, $\rho = 4a$.

Ex. 4. A heavy particle hangs from a point O by a string of length a . It is projected horizontally with a velocity v such that $v^2 = (2 + \sqrt{3}) ag$. Show that the string becomes slack when it has described an angle

$$\cos^{-1} \left(-\frac{1}{\sqrt{3}} \right).$$

CONSTRAINED MOTION

The string becomes slack when $T = 0$.

Hence, from (5), we have

$$3g \cos \theta + \sqrt{3}g = 0$$

or, $\cos \theta = -\frac{1}{\sqrt{3}}$, giving $\theta = \cos^{-1} \left(-\frac{1}{\sqrt{3}} \right)$.

Ex. 5. A bead slides down a rough circular wire, which is in a vertical plane, starting from rest at the end of the horizontal diameter. When it has described an angle θ about the centre, show that the square of the angular velocity is

$$\frac{2g}{a(1+4\mu^2)} \left((1-2\mu^2) \sin \theta + 3\mu(\cos \theta - e^{-2\mu\theta}) \right),$$

where μ is the coefficient of friction and a is the radius of the circle.

[B. H. 1992; C. H. 2001]

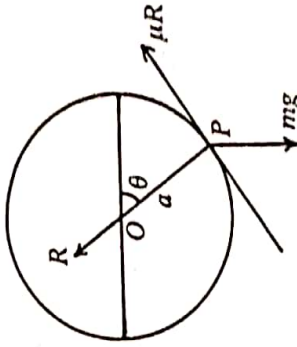


Fig. 13.8 (c)

The forces acting on the particle are

- (i) mg vertically downwards,
- (ii) reaction R normal to the path,
- (iii) friction μR tangentially.

Taking the centre of the circle as pole and the horizontal diameter as initial line, the co-ordinates of P are (a, θ) .

Equations of motion in the tangential and the normal directions are

$$m a \ddot{\theta} = mg \cos \theta - \mu R,$$

$$m a \dot{\theta}^2 = R - mg \sin \theta,$$

since $v = a \dot{\theta}$ and $\rho = a$, so that $\frac{dv}{dt} = a \ddot{\theta}$ and $\frac{v^2}{\rho} = a \dot{\theta}^2$.

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Let the particle be at P at time t such that $AP = s$ and $\angle AOP = \theta$, where A is the initial position of the particle. The forces acting on the particle are

- (i) weight mg acting vertically downwards,
- (ii) tension T acting along the string towards O .

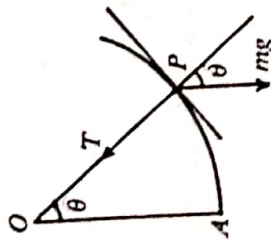


Fig. 13.8 (b)

The equations of motion are

$$m v \frac{dv}{ds} = -mg \sin \theta \quad \dots (1)$$

$$\text{and} \quad m \frac{v^2}{\rho} = T - mg \cos \theta. \quad \dots (2)$$

Now $s = a\theta$, so that $\frac{ds}{d\theta} = a$.

From (1), we have

$$v \frac{dv}{d\theta} = -ag \sin \theta.$$

Integrating, we get

$$v^2 = 2ag \cos \theta + C, \text{ where } C \text{ is a constant.} \quad \dots (3)$$

Initially, at A , $\theta = 0$ and $v^2 = (2 + \sqrt{3}) ag$.

Hence $C = \sqrt{3}ag$.

Substituting in (3), we get

$$v^2 = 2ag \cos \theta + \sqrt{3}ag. \quad \dots (4)$$

From (2) and (4), we have

$$\frac{m}{a} (2ag \cos \theta + \sqrt{3}ag) = T - mg \cos \theta$$

$$\text{or,} \quad T = m (3g \cos \theta + \sqrt{3}g). \quad \dots (5)$$

CONSTRAINED MOTION

Eliminating R, we get

$$v \frac{dv}{ds} - \mu \frac{v^2}{\rho} = g (\mu \cos \psi - \sin \psi)$$

or, $\frac{1}{2} \frac{d}{ds} (v^2) - \mu \frac{v^2}{\rho} = g (\mu \cos \psi - \sin \psi)$

or, $\frac{dv^2}{d\psi} - 2\mu v^2 = 2g\rho (\mu \cos \psi - \sin \psi)$.

Integrating, we get

$$v^2 e^{-2\mu\psi} = 2g \int \rho e^{-2\mu\psi} (\mu \cos \psi - \sin \psi) d\psi.$$

Now, for a cycloid, $s = 4a \sin \psi$ and $\rho = 4a \cos \psi$.

Hence $v^2 e^{-2\mu\psi} = 8ag \int e^{-2\mu\psi} \cos \psi (\mu \cos \psi - \sin \psi) d\psi.$

Put $z = e^{-\mu\psi} (\mu \cos \psi - \sin \psi)$.

Then $dz = -e^{-\mu\psi} (1 + \mu^2) \cos \psi d\psi$.

Hence we have

$$v^2 e^{-2\mu\psi} = -\frac{8ag}{1 + \mu^2} \int z dz = -\frac{4ag}{1 + \mu^2} z^2 + A,$$

A being a constant

$$= -\frac{4ag}{1 + \mu^2} (\mu \cos \psi - \sin \psi)^2 e^{-2\mu\psi} + A. \dots (1)$$

(i) In this case, the conditions are $v = 0$ when $\psi = \frac{1}{2}\pi$ and $\psi = 0$.

Putting these values in (1), we get

$$A = \frac{4ag}{1 + \mu^2} e^{-\pi\mu} = \frac{4ag}{1 + \mu^2} \mu^2.$$

Therefore $\mu^2 = e^{-\pi\mu}$, whence $\mu^2 e^{\pi\mu} = 1$.

(ii) In this case, the conditions are $v = 0$ when $\psi = \theta$ and $\psi = 0$.
Putting these values in (1), we get

$$A = \frac{4ag}{1 + \mu^2} (\mu \cos \theta - \sin \theta)^2 e^{-2\mu\theta} = \frac{4ag}{1 + \mu^2} \mu^2.$$

Therefore $\mu^2 e^{2\mu\theta} = (\sin \theta - \mu \cos \theta)^2$

or, $\mu e^{\mu\theta} = \sin \theta - \mu \cos \theta$.

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Eliminating R from the equations, we get

$$\ddot{\theta} + \mu \dot{\theta}^2 = \frac{g}{a} (\cos \theta - \mu \sin \theta)$$

or, $\frac{d}{d\theta} (\dot{\theta}^2) + 2\mu \dot{\theta}^2 = \frac{2g}{a} (\cos \theta - \mu \sin \theta)$.

Integrating, we get

$$\dot{\theta}^2 e^{2\mu\theta} = \frac{2g}{a} \int e^{2\mu\theta} (\cos \theta - \mu \sin \theta) d\theta$$

$$= \frac{2g}{a} \left(\frac{2\mu e^{2\mu\theta} \cos \theta + e^{2\mu\theta} \sin \theta}{1 + 4\mu^2} - \mu \frac{2\mu e^{2\mu\theta} \sin \theta - e^{2\mu\theta} \cos \theta}{1 + 4\mu^2} \right) + C$$

$$= \frac{2g}{a} \frac{e^{2\mu\theta}}{1 + 4\mu^2} (3\mu \cos \theta + (1 - 2\mu^2) \sin \theta) + C.$$

Initially, $\theta = 0 = \dot{\theta}$. Hence $\frac{2g}{a} \cdot \frac{3\mu}{1 + 4\mu^2} + C = 0$.

Therefore $C = -\frac{2g}{a} \frac{3\mu}{1 + 4\mu^2}$.

Hence the result follows.

Ex. 6. A heavy particle slides down a rough cycloid of which the coefficient of friction is μ . Its base is horizontal and vertex downwards. Show that

(i) if it starts from rest at the cusp and comes to rest at the vertex, then $\mu^2 e^{\mu\pi} = 1$; [C. H. 1986 ; B. H. 1990]

(ii) if it starts from rest at a point where the tangent makes an angle θ with the horizontal and comes to rest at the vertex, then $\mu e^{\mu\theta} = \sin \theta - \mu \cos \theta$. [B. H. 1984 ; C. H. 1984]

The motion being down the cycloid, the friction μR acts up the tangent.

The equations of motion are then

$$m \frac{dv}{ds} = \mu R - mg \sin \psi$$

and $m \frac{v^2}{\rho} = R - mg \cos \psi$.